# Ionisation by magnetic reconnection events in T Tauri discs



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#### I. Context

#### **Regions of interest**



Interaction region between star magnetic field and disc magnetic field

#### **Reconnection regions**



**Energetic particles** are produced that will **ionise** the inner disc

#### Interests of ionisation in protoplanetary discs

#### Source of **heating** of disc and jet Initiate disc (prebiotic) **chemistry** Controls **accretion**

The study : **Ionisation** rates in T Tauri discs due to magnetic reconnection events

## Procedure

$$\zeta(N) = 2\pi \int_{E_{ion}}^{\infty} j(E,N) \sigma_{ion,k}(E) dE \quad (s^{-1})$$

j(E, N): the propagated spectrum

$$j(E) = j_0(E_0) \frac{L(E_0)}{L(E)}$$

Structure : I. Disc Model II. Injection Model III.Particle Disc interaction

**IV.Results** 

Injection model

Particle – Disc interaction

Padovani et al. 2009

## II. Protoplanetary Disc Model



## Chemical Model: ProDiMo (Woitke P. 2009)



Composition: Plasma, H, H<sub>2</sub> and He



# Flare append at the **disc edge**

#### **Magnetic Configurations**



Vertical : Référence

Hyperbolic : Disque standard (Blandford, R. D., & Payne, (1982))

**Quartic** : Differential accretion in the disc

#### **Toroidal Magnetic Configurations**



Particles explore **different** column densities for each configuration

## III. Particle and Radiation Model



#### Density as a function of flare size



Getman, K. et al. (2008) The Astrophysical Journal.

Bremsstrahlung Luminosity (erg/s):

$$L_X = 1.4 \times 10^{26} L_{10}(T)^3 n_{e,10}(T)^2 T_6^{1/2} g_E$$

 $L_{10}$  : size of the flare over  $10^{10}$  cm  $n_{e,10}$  : electron density over  $10^{10}$  cm<sup>-3</sup>  $T_6$  : temperature over  $10^6$  K  $g_B$ : Bolometric Gaunt factor

## ProDiMo computes the chemical structure of the disc for different flare temperatures



Atomic hydrogen is present deeper in the disc for hot flares with strong X-ray emission

#### **Particle Injection**



**Power law** spectrum as in solar flares

$$j_0(E,T_F) \sim \left(\frac{E}{E_{inj}(T_F)}\right)^{-\delta} \exp\left(-\frac{E}{100 \ MeV}\right)$$

 $n_{NTh}$  : density of non-thermal particles  $E_{inj}$ : injection energy  $T_F$ : temperature of the flare

## IV. Disc Ionisation Model

#### The loss function

$$\frac{dE}{dN} = -L(E)$$

The loss function describes a maximum of **interaction processes** 

#### **Loss functions**



Loss function depends on the **particle** and the **medium** crossed

$$\overline{L}(E,s) = \sum_{i} f_{i}(s) L_{i}(E) \quad i = H, H_{2}, He$$

$$f_{i} = \frac{1}{s} \int_{0}^{s} \frac{n_{i}(s')}{n_{tot}(s')} ds'$$

We build a mean loss function at **each position** 

$$j(E,N) = j_0(E_0) \frac{\overline{L}(E_0)}{\overline{L}(E)}$$

Padovani, M. et al (2009) Astronomy & Astrophysics.

Continuous slowing down approximation gives the propagated spectrum

$$\zeta(N) = 2\pi \int j(E, N) (1 + \phi(E)) \sigma_{ion}(E) dE$$

 $\sigma_{ion}$ : ionisation cross section  $\phi(E)$ : ionisation by secondary particles lonisation rates
obtained from
propagated
spectrum

## V. Results





Hotter flare, higher ionisation rate

Even weak flares are a dominant source of ionisation



Lower index, higher ionisation rate

 $\begin{array}{l} \textbf{Dominant source} \\ \textbf{of ionisation for} \\ \boldsymbol{\delta} < \textbf{5} \end{array}$ 

#### **Toward a more Predictive Model**

Ionisation rate are overestimated due to very localised results



Need of a spatial and time averaged model