

# Turbulent solar and stellar winds



V. Réville, M. Velli, N. Fargette, A. Rouillard, B. Lavraud, S. Parenti, S. Brun, A. Strugarek, M. Shoda, PSP and SolO teams.



# Atmospheres of solar-like stars

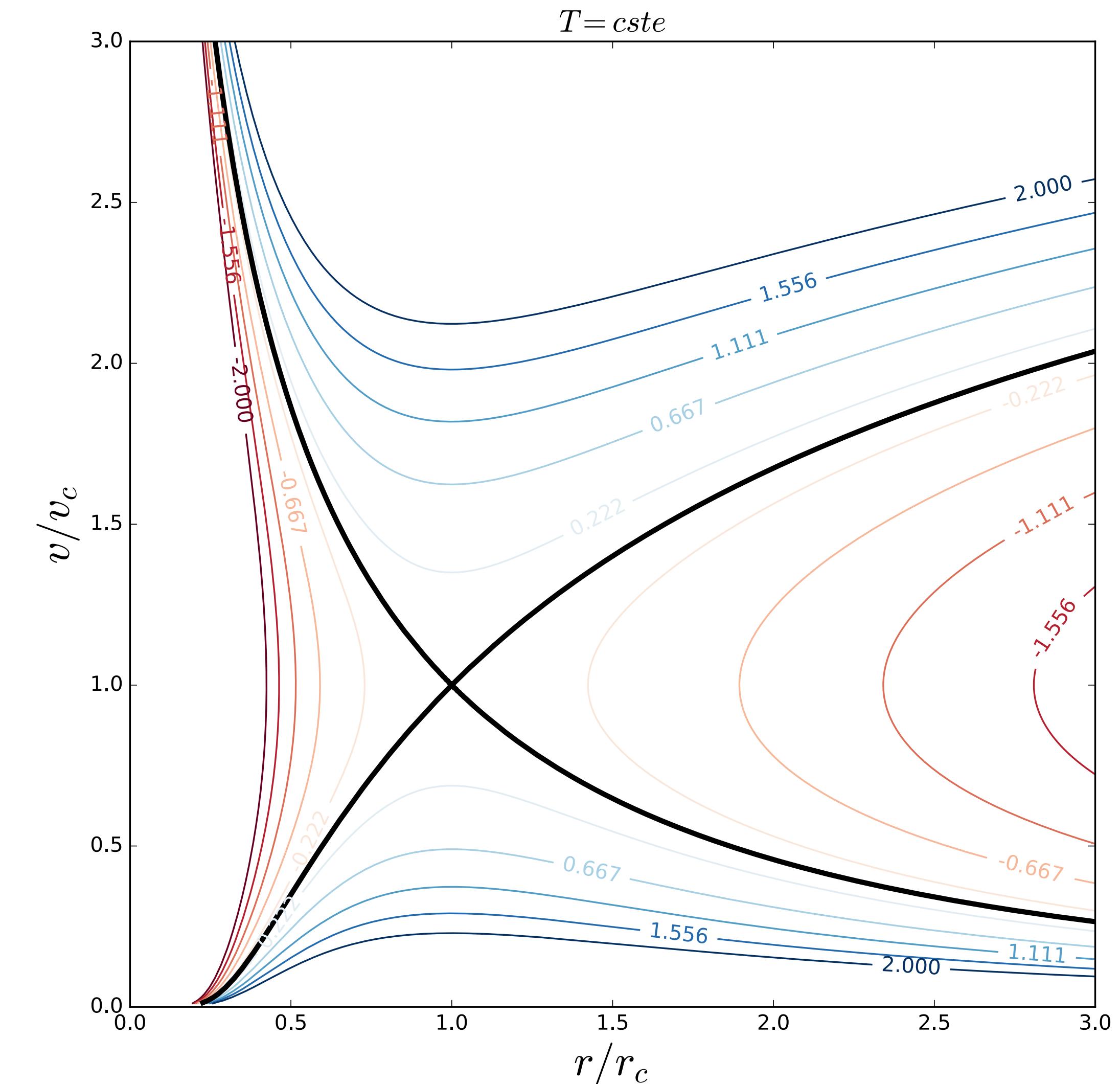
## Wind = MK corona under gravity potential

- Hydrogen plasma of 1 MK
- Spherical geometry with a central object
- Very fast thermal conduction (electrons) ~isothermal

$$r_c = \frac{GM_\star}{2c_s^2}$$

[Parker 1958]

$$v_c = c_s = \sqrt{\frac{\partial p}{\partial \rho}}$$

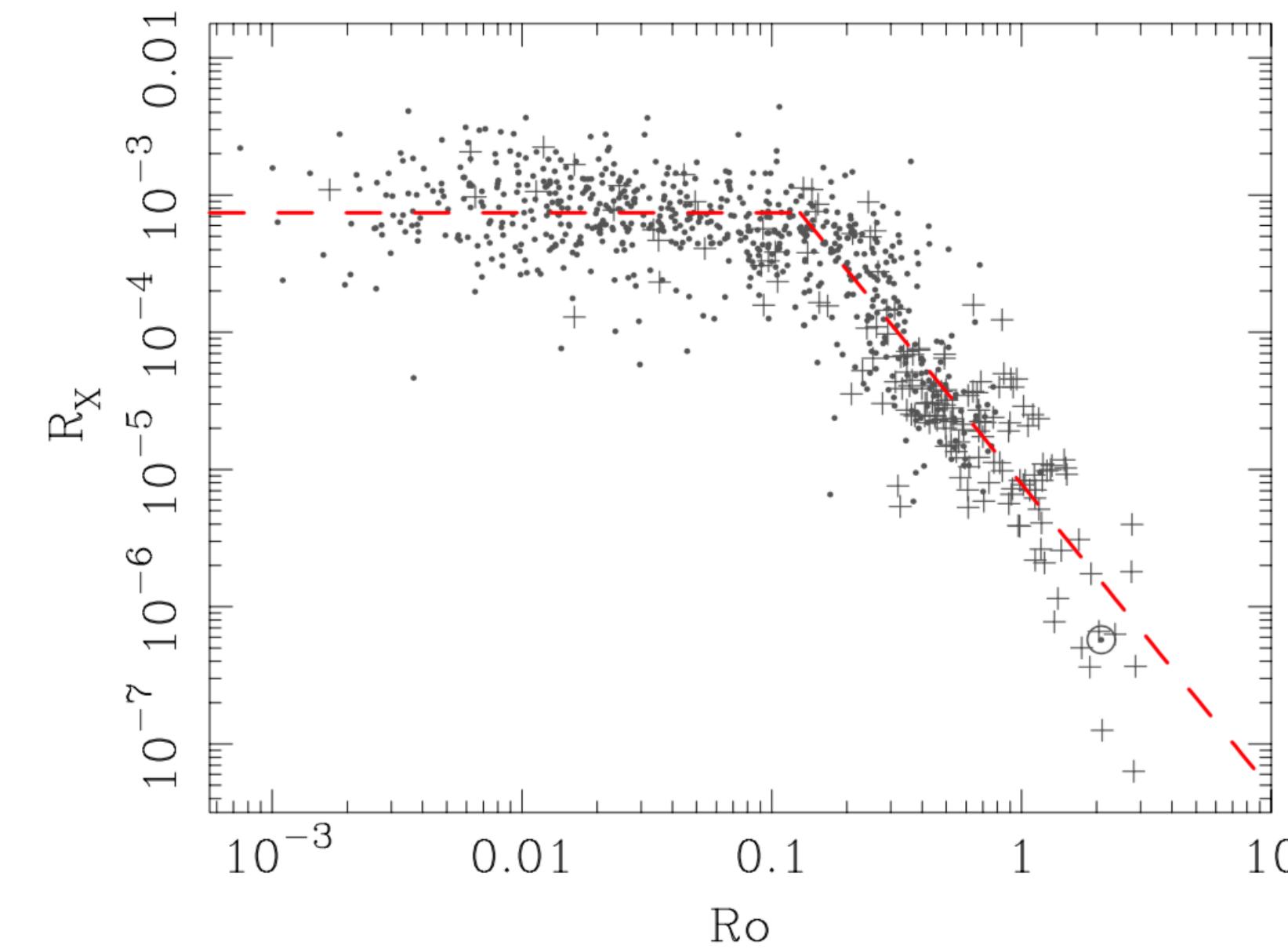


# Stellar winds: observables

## Transitioning to active/fast rotating stars

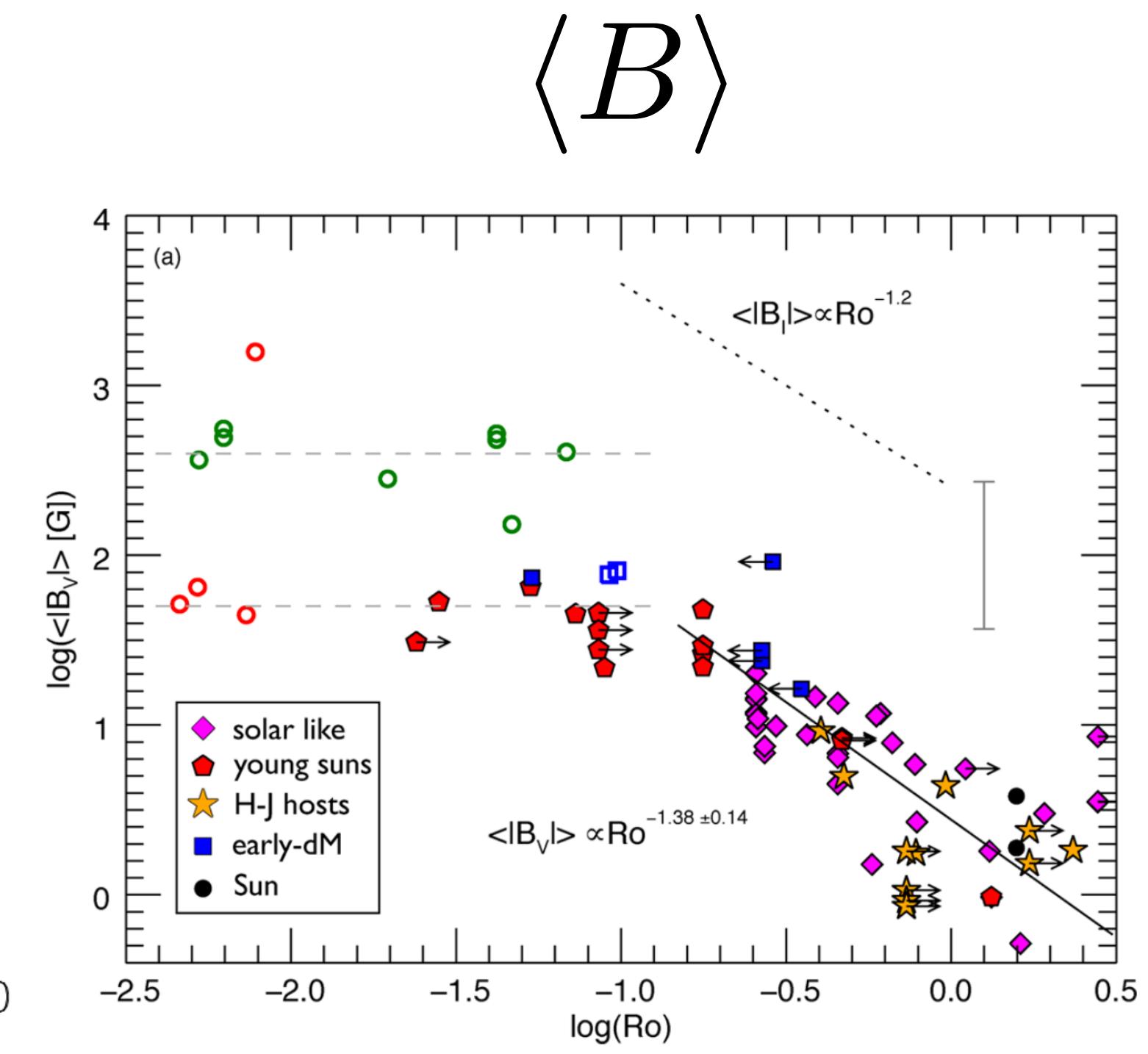
$$Ro = \frac{1}{\Omega_\star \tau_{\text{cz}}}$$

Rotation



[Benbakoura, Réville et al. 2021]

X-ray



[Vidotto et al. 2014]

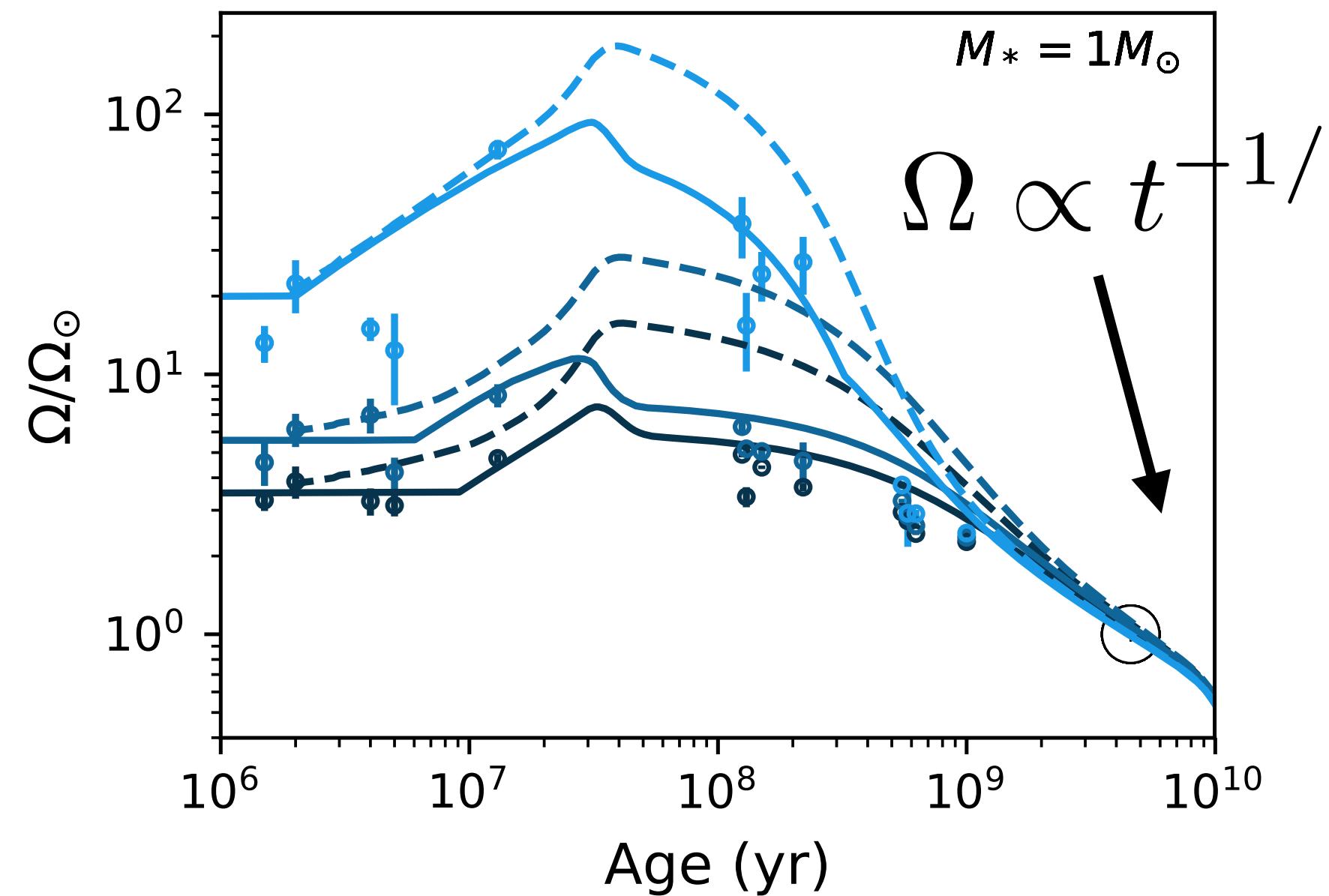
- Saturation phase for active stars in braking, X-ray and magnetic field

# Stellar winds: observables

## Transitioning to active/fast rotating stars

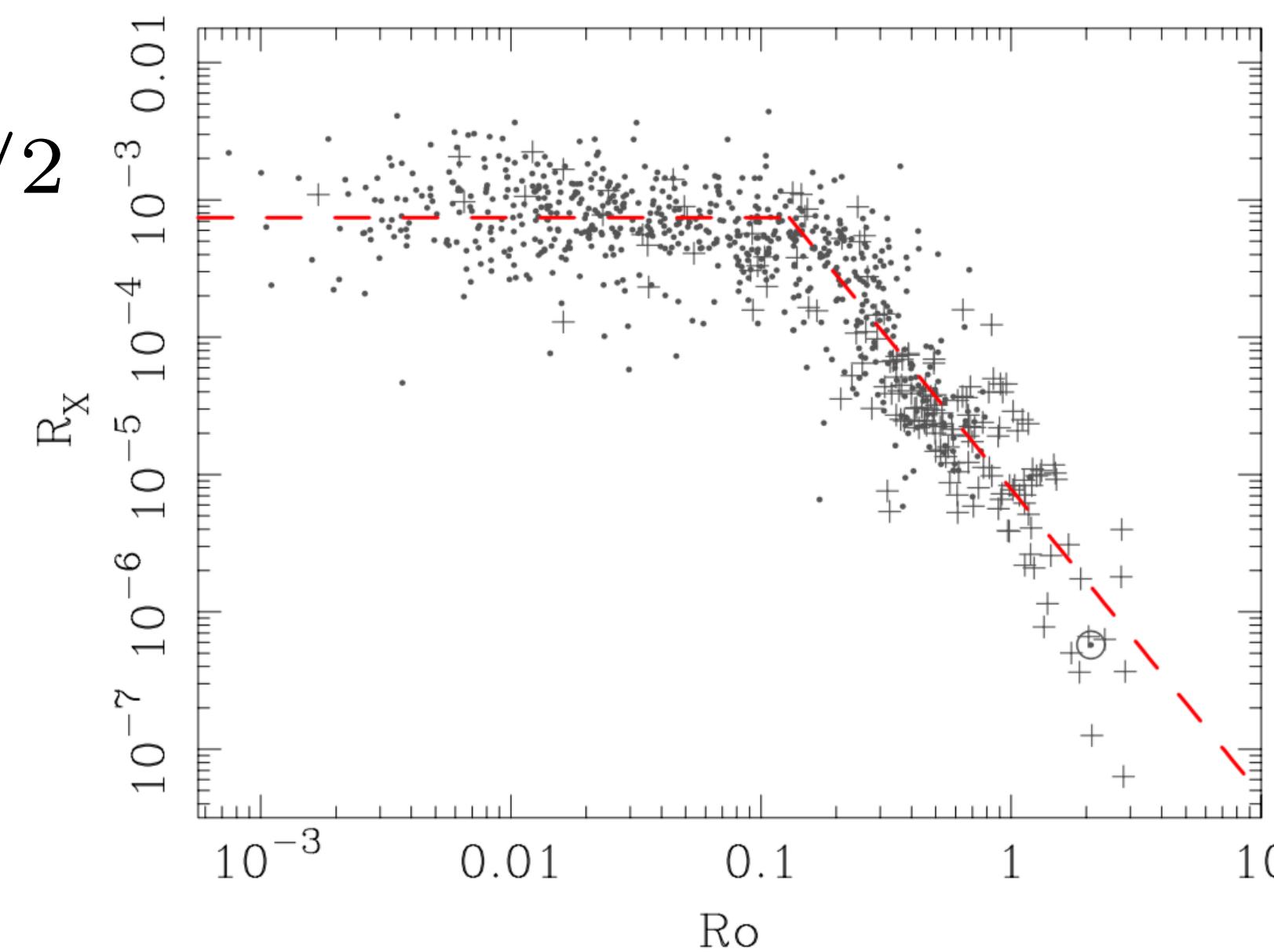
$$Ro = \frac{1}{\Omega_\star \tau_{\text{cz}}}$$

Rotation



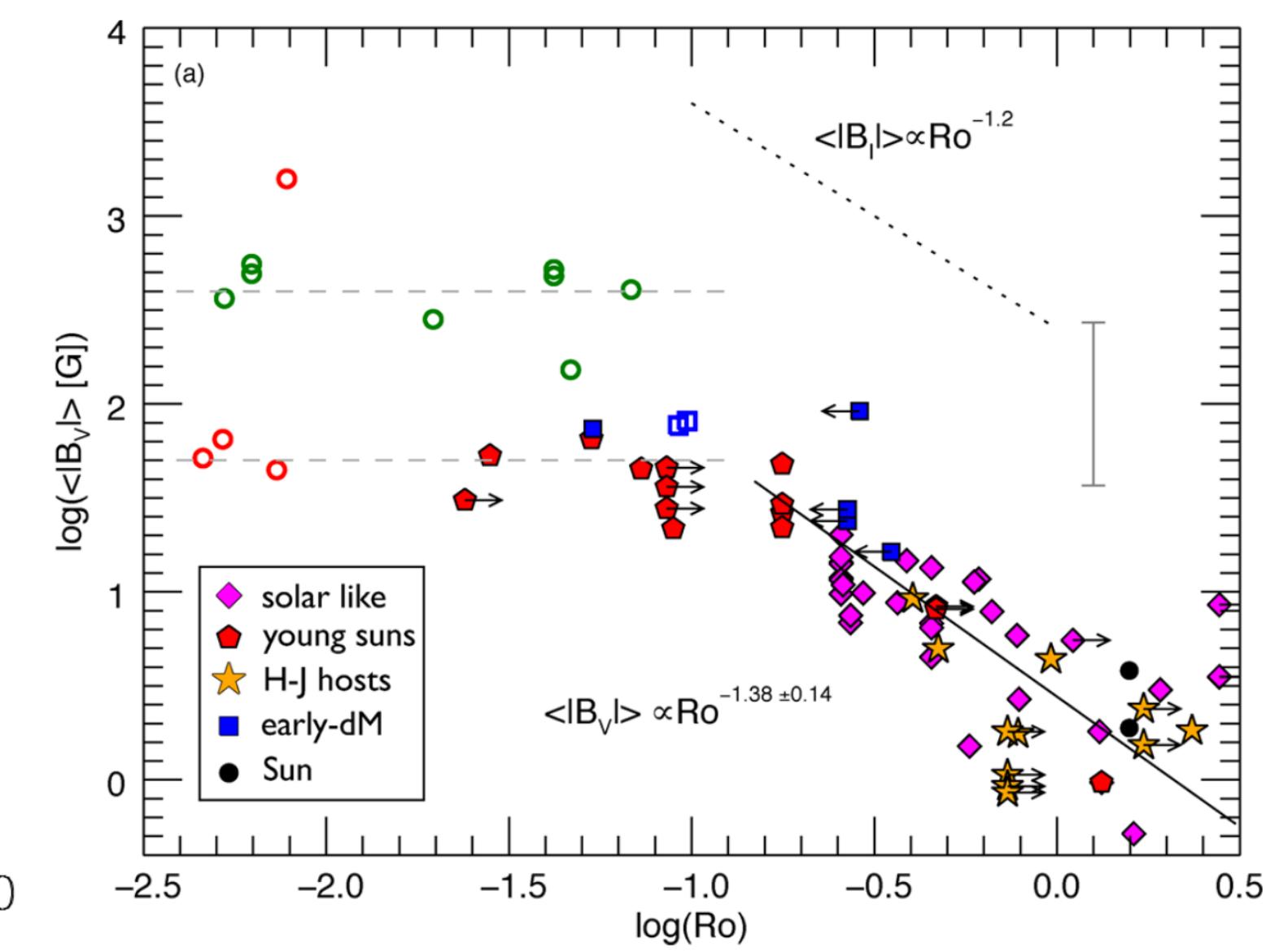
[Benbakoura, Réville et al. 2021]

X-ray



[Wright et al. 2011]

$\langle B \rangle$



[Vidotto et al. 2014]

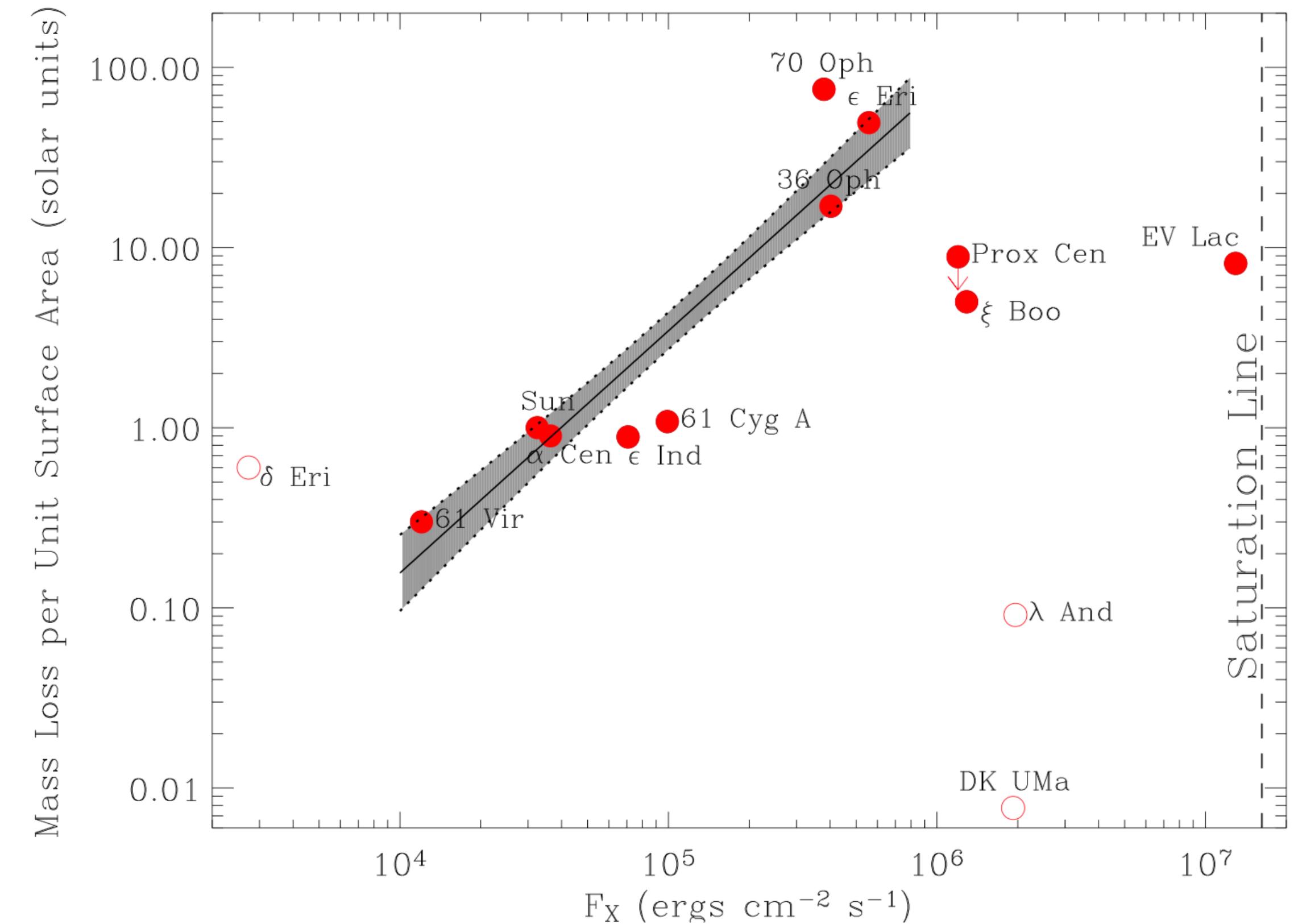
- Saturation phase for active stars in braking, X-ray and magnetic field

# Stellar winds: observables

## Mass loss saturation

- Very little (indirect) mass loss observations
- Correlation with  $F_X \rightarrow$  Rotation, Mag field
- Saturation (before X-ray sat.)
- Largely unexplained

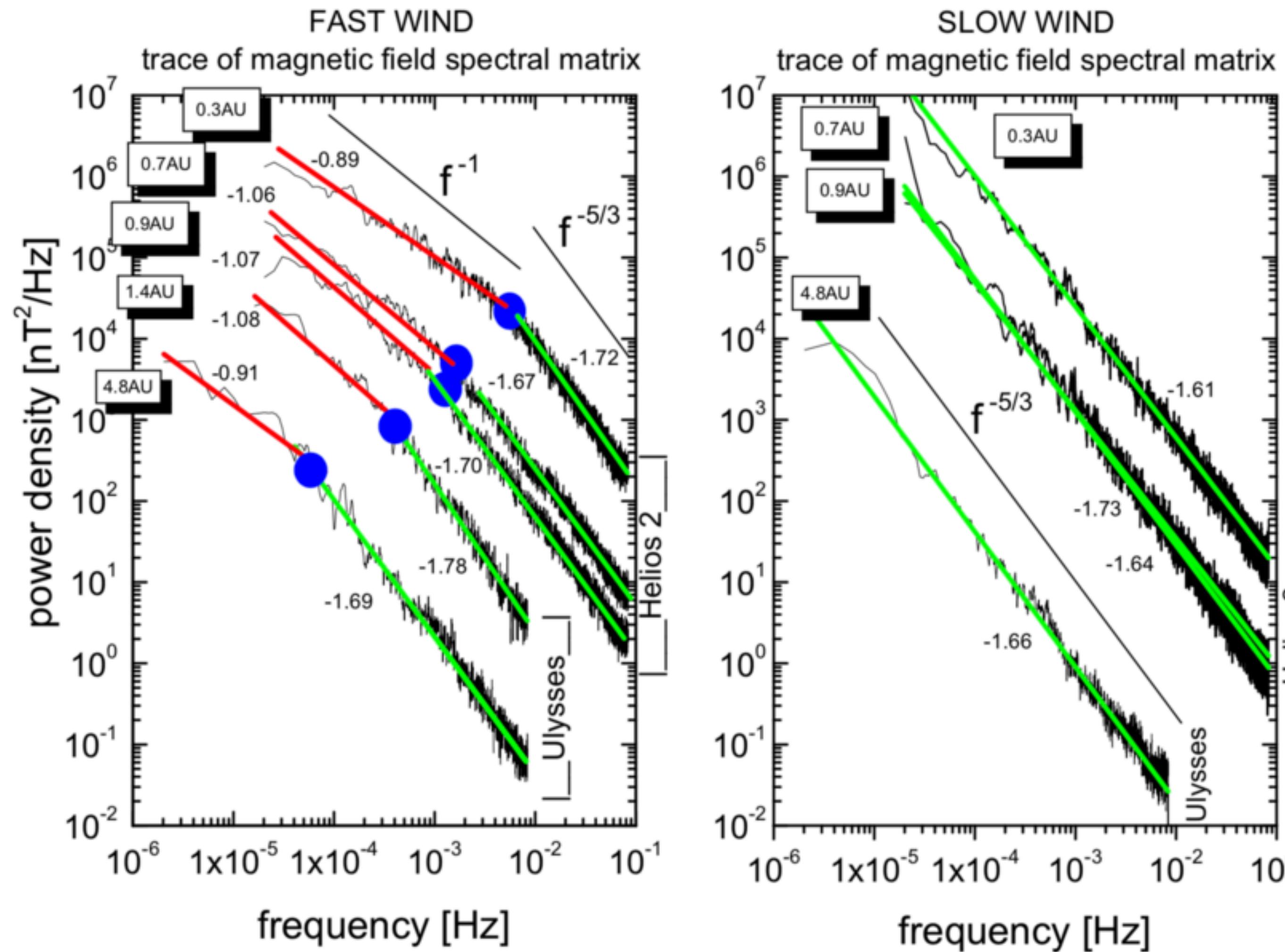
[Wood et al. 2005]



# **Turbulence as the main process of solar wind acceleration and coronal heating**

# Turbulence in the solar wind

## Observations



- The solar wind is a large Reynolds (Lundquist) number system ( $> 10^{12}$ ).
- Turbulence is the universal way to transfer energy from large scales to small, kinetic scales (heating).
- Observed in the solar wind with typical Kolmogorov spectra  $f^{-5/3}$ .

[Bruno & Carbone 2013]

# Alfvén wave turbulence 3D MHD model

## Equations

$$\partial_t \rho + \nabla \cdot [\rho \mathbf{v}] = 0,$$

$$\partial_t (\rho \mathbf{v}) + \nabla \cdot [\rho \mathbf{v} \mathbf{v} - \mathbf{B} \mathbf{B} + \mathbf{I}(p + \mathcal{E}/2)] = -\rho \nabla \Phi,$$

$$\partial_t (E + \mathcal{E} + \rho \Phi) + \nabla \cdot [(E + p + \mathcal{E}/2 + \rho \Phi) \mathbf{v} - \mathbf{B}(\mathbf{v} \cdot \mathbf{B}) + \mathbf{v}_g^+ \mathcal{E}^+ + \mathbf{v}_g^- \mathcal{E}^-] = Q,$$

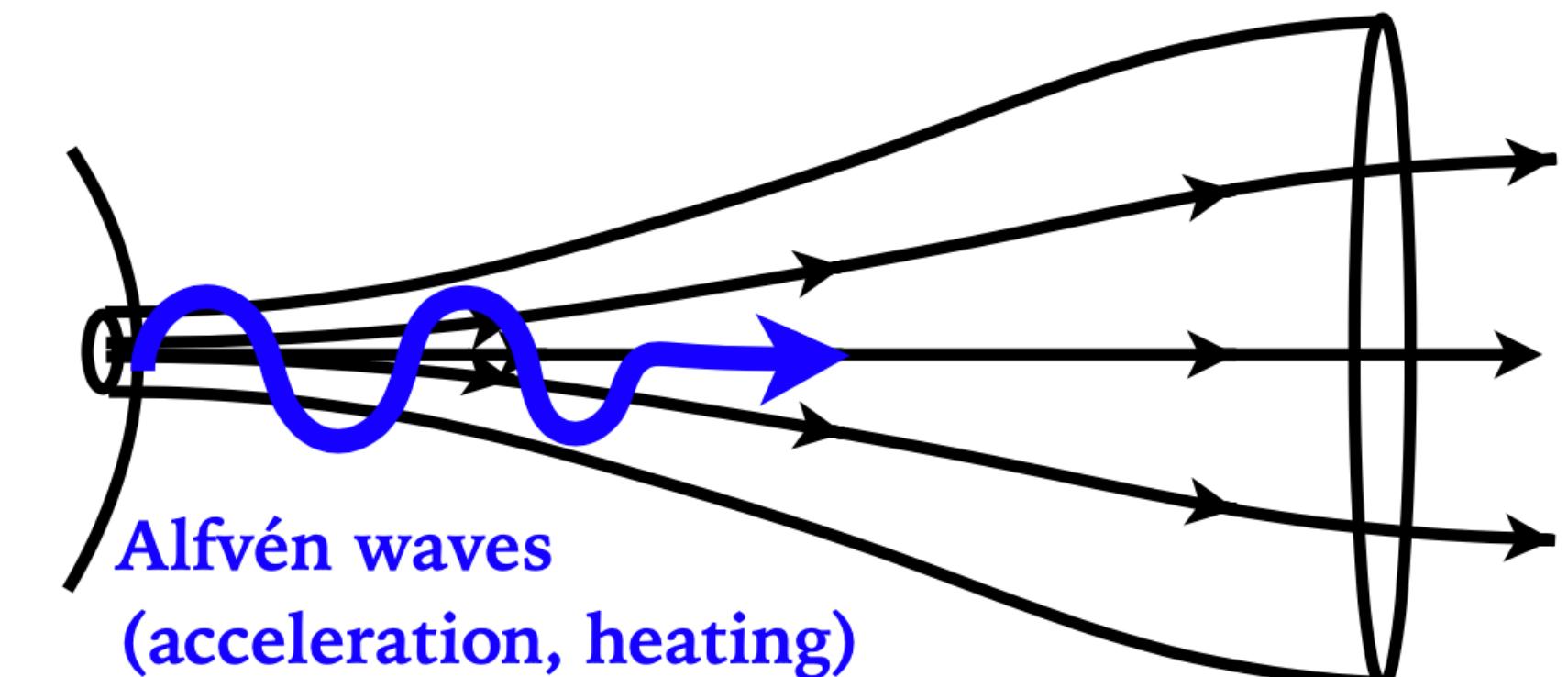
$$\partial_t \mathbf{B} + \nabla \cdot [\mathbf{v} \mathbf{B} - \mathbf{B} \mathbf{v}] = \eta \nabla \times \mathbf{B},$$

$$\partial_t \mathcal{E}^\pm + \nabla \cdot [(\mathbf{v} \pm \mathbf{v}_A) \mathcal{E}^\pm] = -\frac{\mathcal{E}^\pm}{2} \nabla \cdot \mathbf{v} - Q_w^\pm,$$

- Core = PLUTO code (open source) [Mignone et al. 2007,2012]
- Physics of Alfvén wave propagation and dissipation [Réville et al. 2020]

$$Q = Q_h + Q_w - Q_c - Q_r$$

↑  
 Volume heating  
 Wave heating  
 Conduction  
 Radiation



$$\mathcal{E}^\pm = \rho \frac{|z^\pm|^2}{4}$$

$$Q_w^\pm = \rho \frac{|z^\pm|^2}{8\lambda_c} (\mathcal{R}|z^\pm| + |z^\mp|)$$

Reflection coeff.

$\mathcal{R} = 0.1$

# Parker Solar Probe

## A quick presentation of the mission

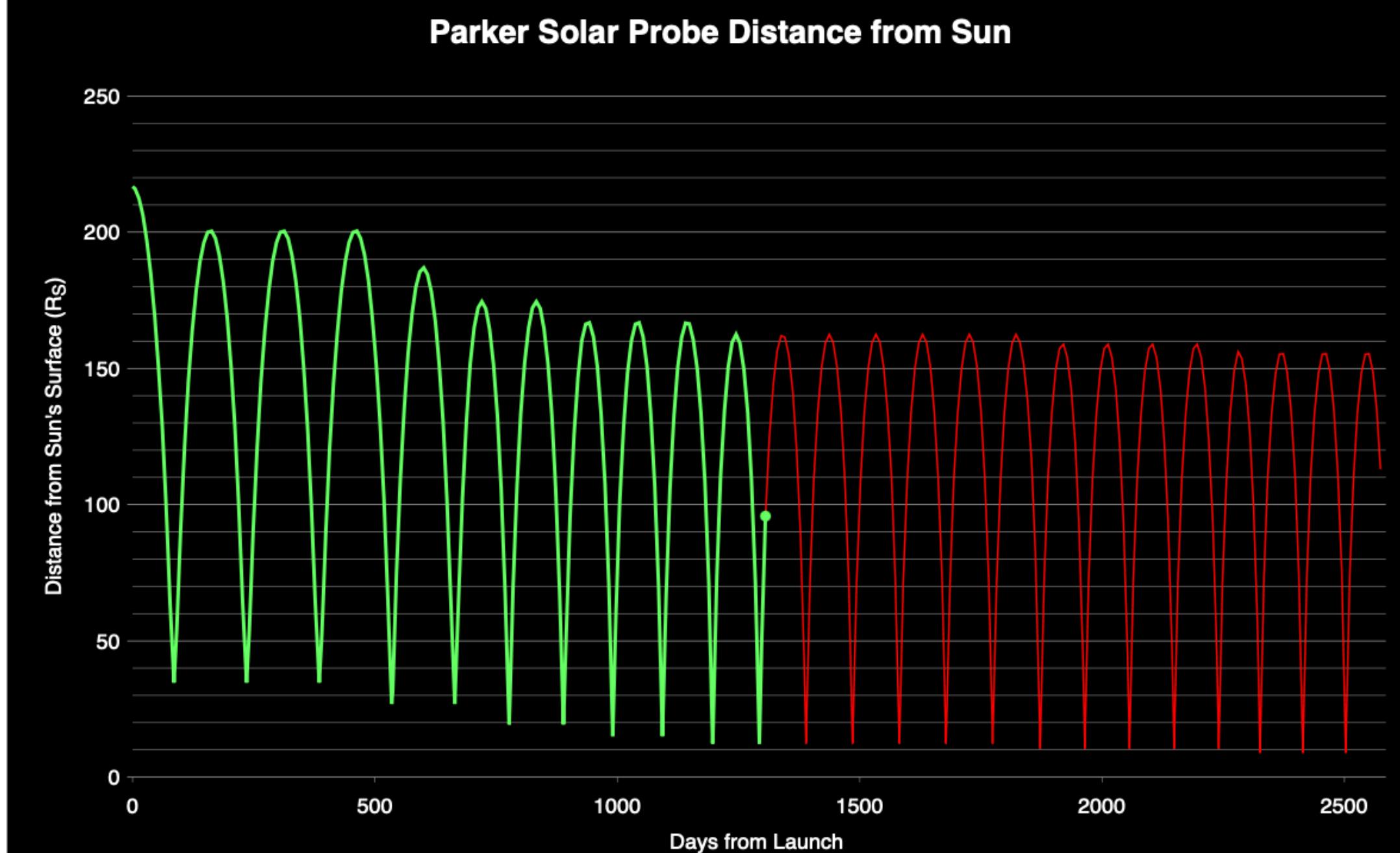
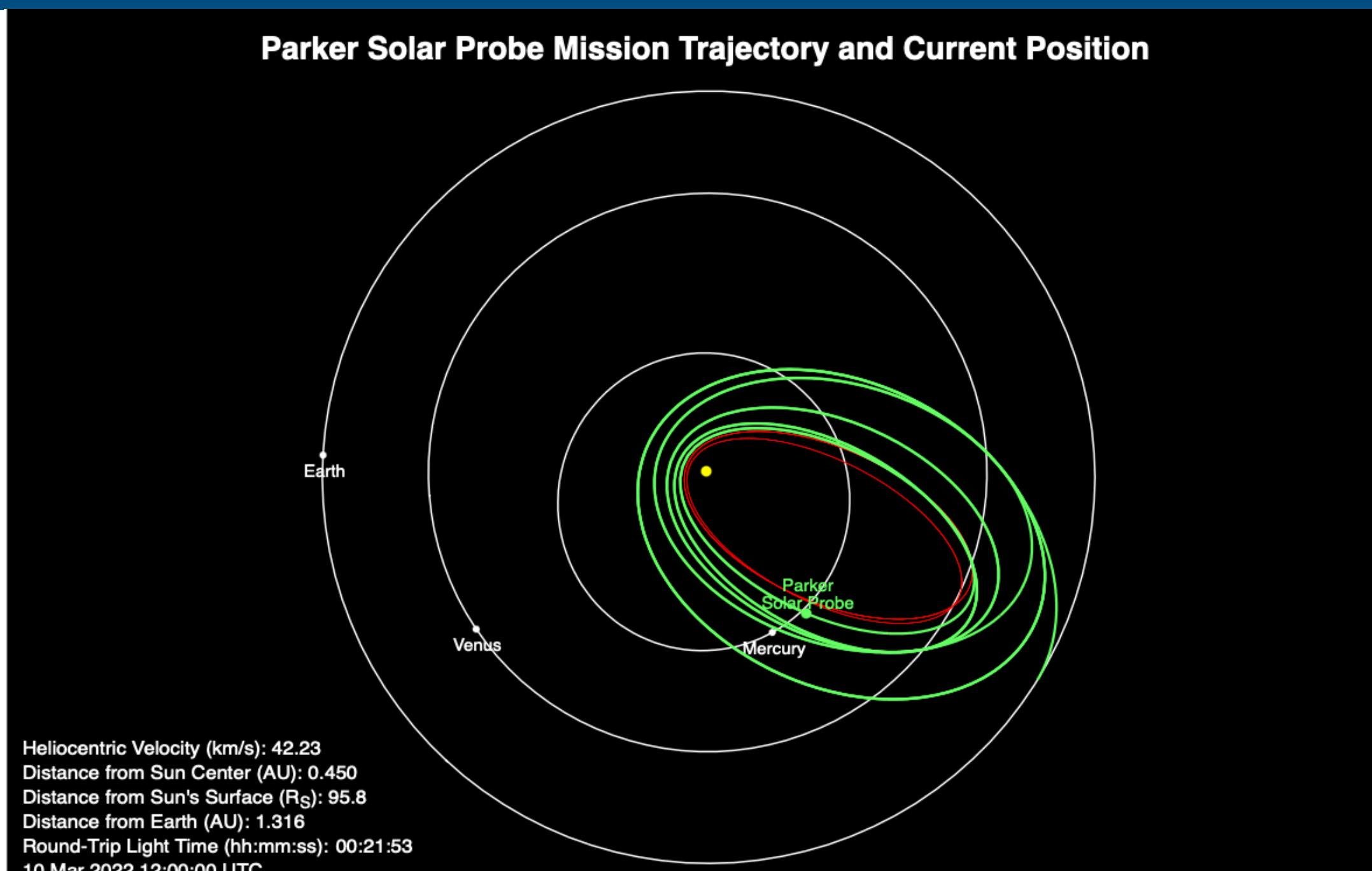
- Launched August 2018
- 4 instruments suite (FIELDS, SWEAP, ISOIS, WISPR)
- First orbit was already the closest we'd been to the Sun
- Closest distance below 10 Sun, i.e., below the Alfvén surface
- Unravel the turbulence and heating mechanisms of the solar wind



# Parker Solar Probe

## A quick presentation of the mission

- Launched August 2018
- 4 instruments suite (FIELDS, SWEAP, ISOIS, WISPR)
- First orbit was already the closest we'd been to the Sun
- Closest distance below 10 Sun, i.e., below the Alfvén surface
- Unravel the turbulence and heating mechanisms of the solar wind

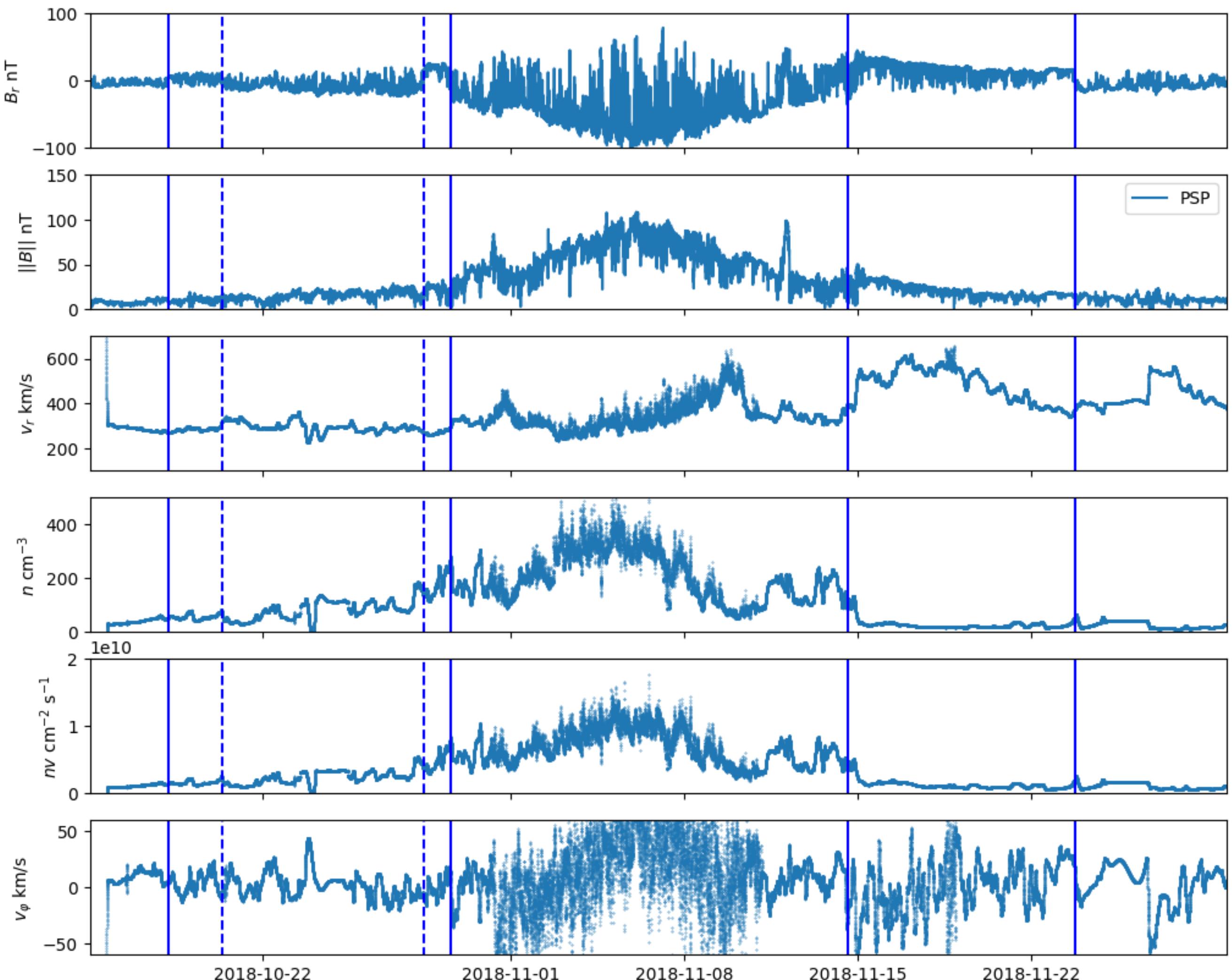


# Alfvén wave turbulence 3D MHD model

## Comparisons with PSP E1 data

[Réville et al. 2020]

- Particles (SWEP) and magnetic field data (FIELDS)
- Single observational input : the magnetic field map (ADAPT) the day of perihelion.
- Average structures, change of polarity are very well reproduced.
- Fast evolving structure (switchbacks) are not meant to be reproduced.

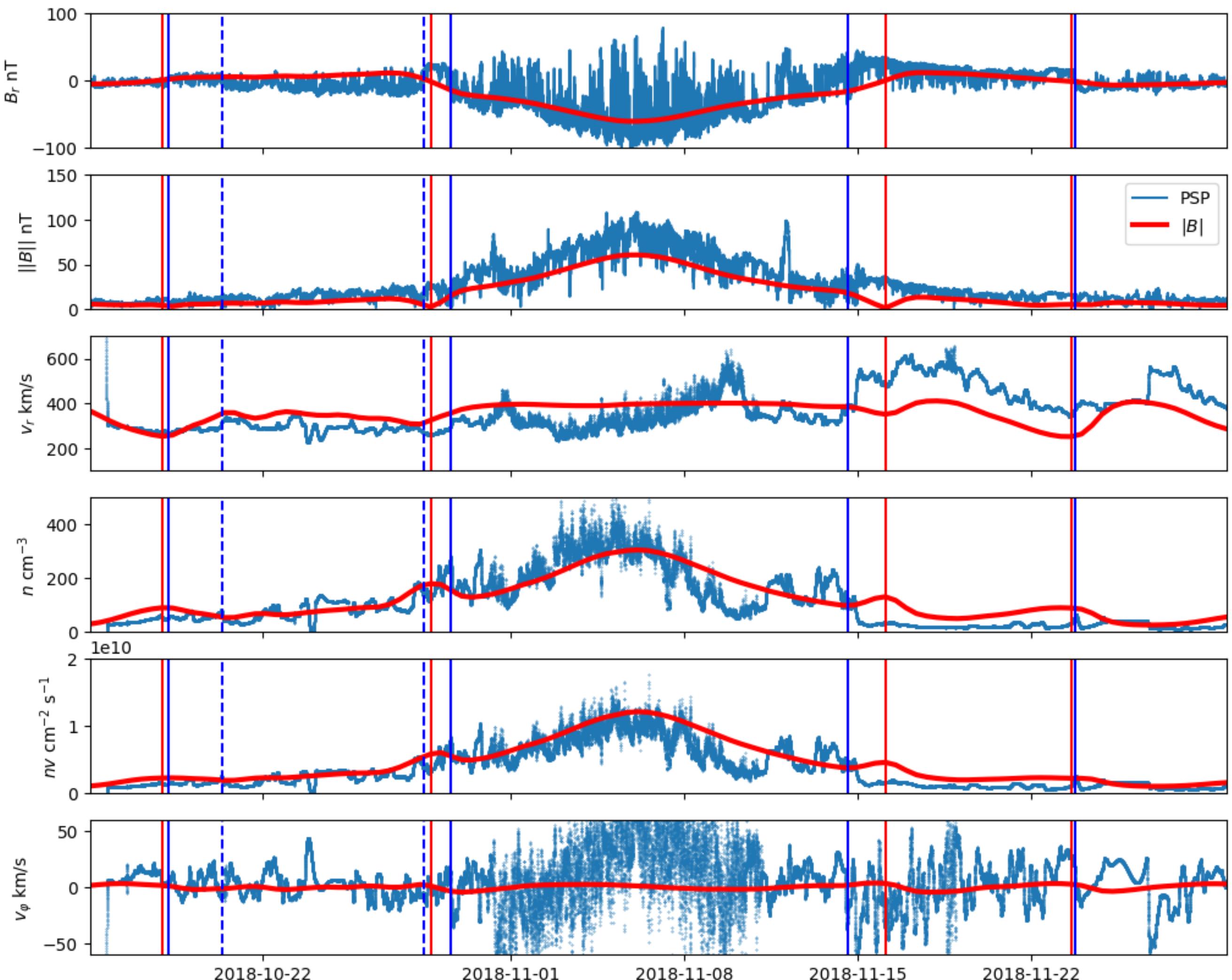


# Alfvén wave turbulence 3D MHD model

## Comparisons with PSP E1 data

[Réville et al. 2020]

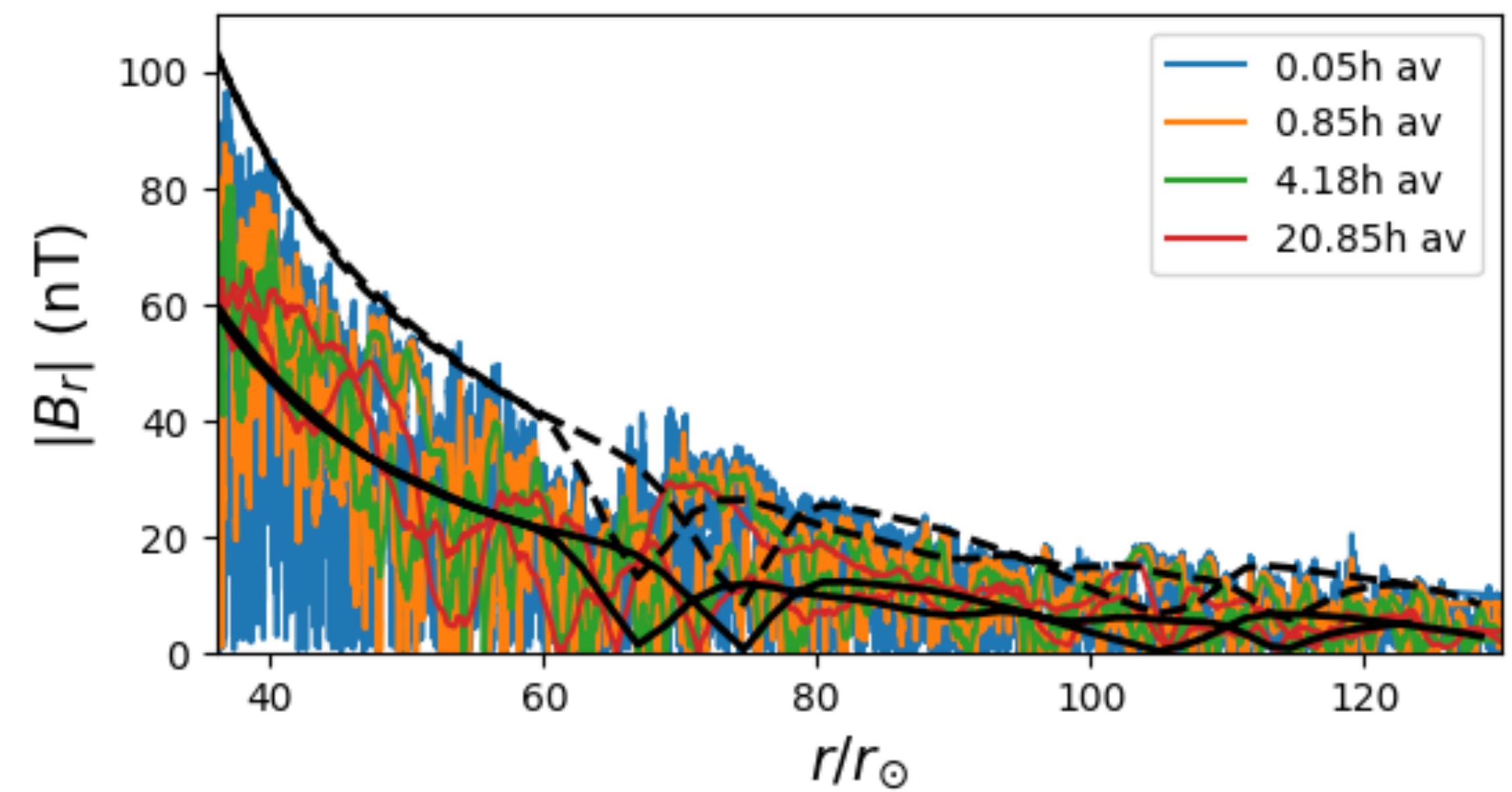
- Particles (SWEP) and magnetic field data (FIELDS)
- Single observational input : the magnetic field map (ADAPT) the day of perihelion.
- Average structures, change of polarity are very well reproduced.
- Fast evolving structure (switchbacks) are not meant to be reproduced.



# Alfvén wave turbulence 3D MHD model

## Perturbations and switchbacks

- Averaging perturbations we fall back on the field amplitude of the model.
- The amplitude of the perturbations in the model is consistent with the observations.
- It proves that AW turbulence models are very good at reproducing solar wind properties.
- Switchbacks are 3D non linear Alfvén waves, probably contributing to the turbulence.



$$\delta b^\pm = \sqrt{\mu_0 \mathcal{E}^\pm}$$

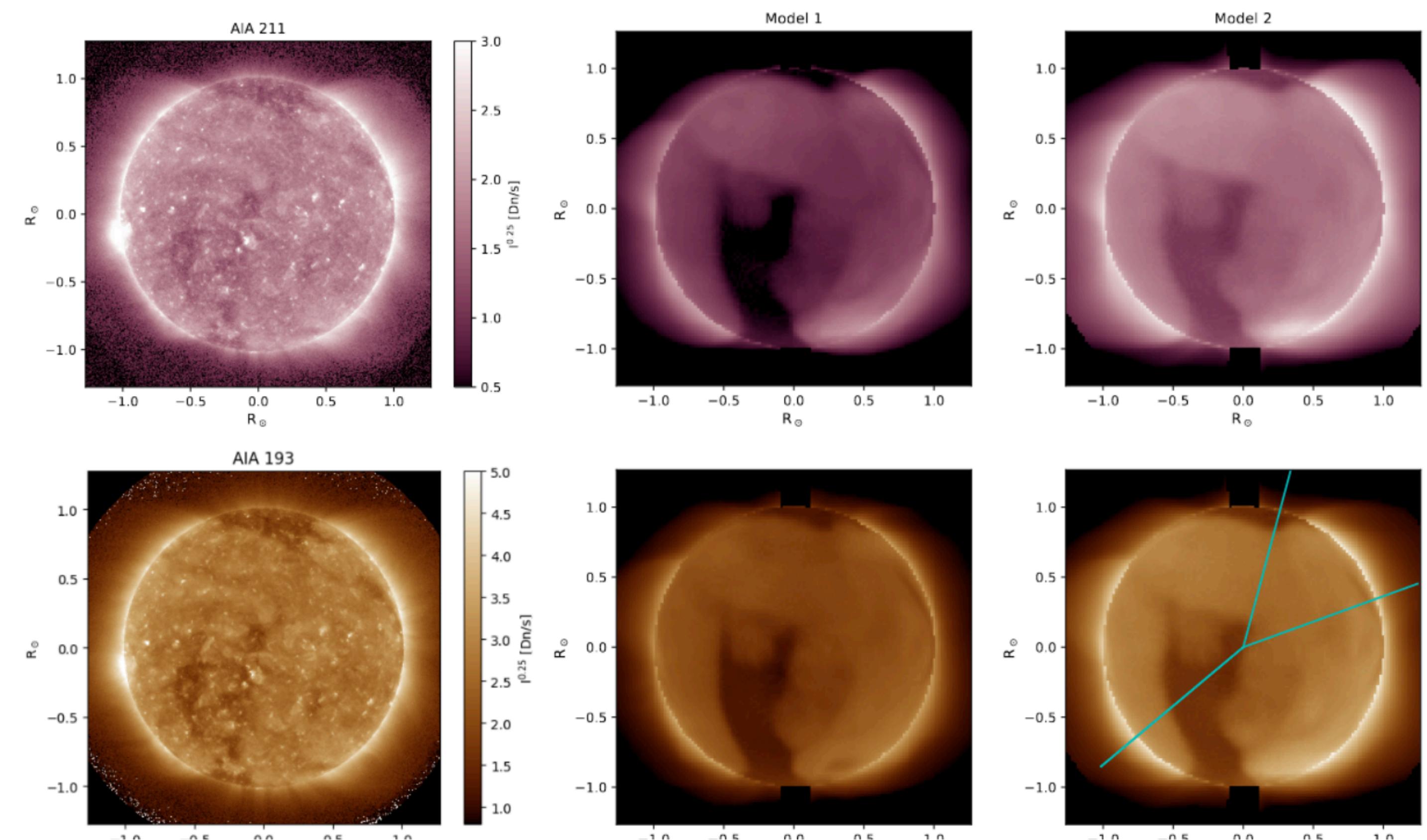
# Alfvén wave turbulence 3D MHD model

## Synthetic remote observations

[Parenti, Réville et al. 2022, ApJ]

- EUV instruments image the solar atmosphere using lines from strongly ionized ions (e.g. Fe)
- SDO/AIA probes temperatures ranging from 0.5 to 2-3 MK.
- We use the instrument response to compute the synthetic emissions from the model

$$I = \int_{LOS} n^2 \mathcal{R}(n, T) dl$$

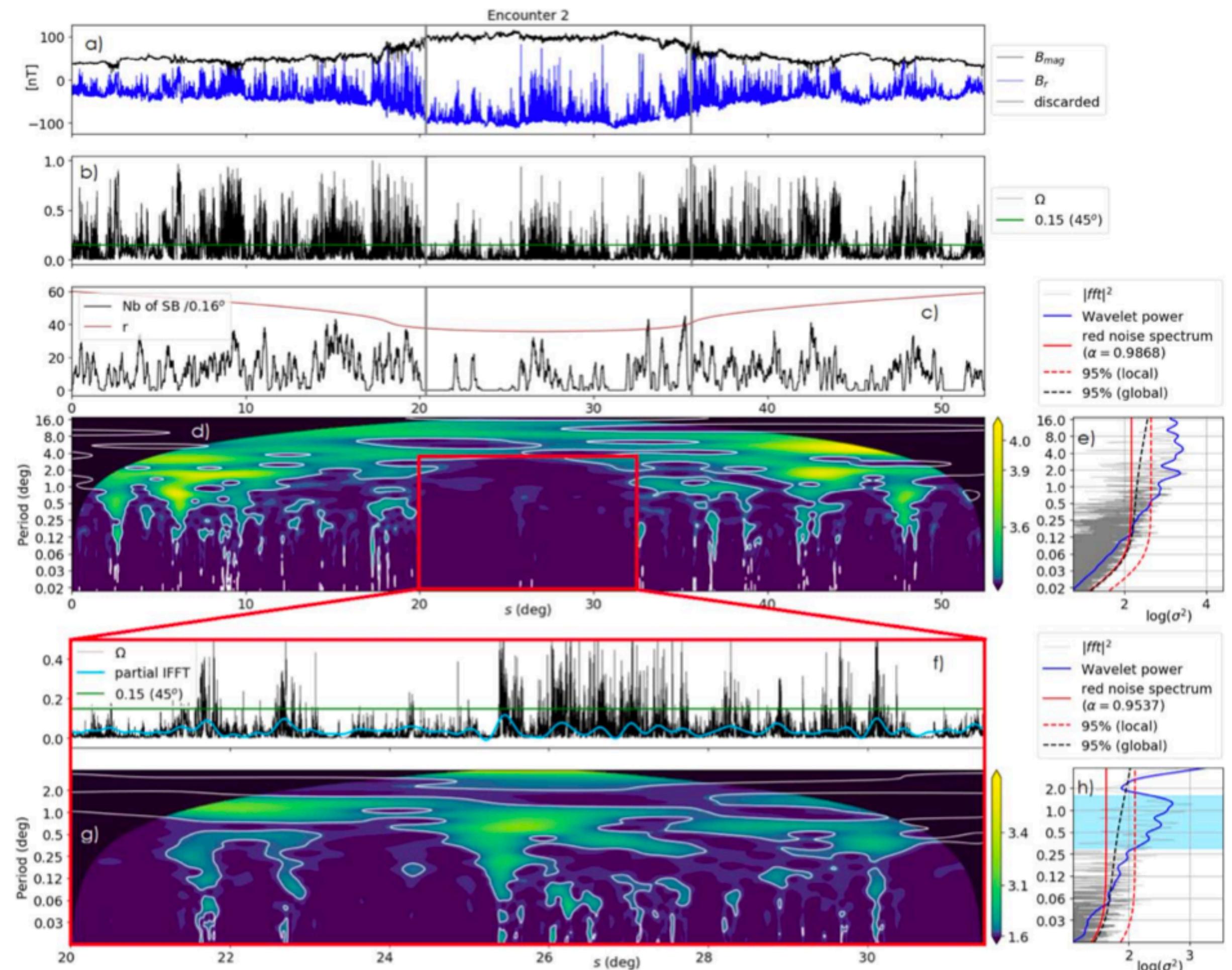
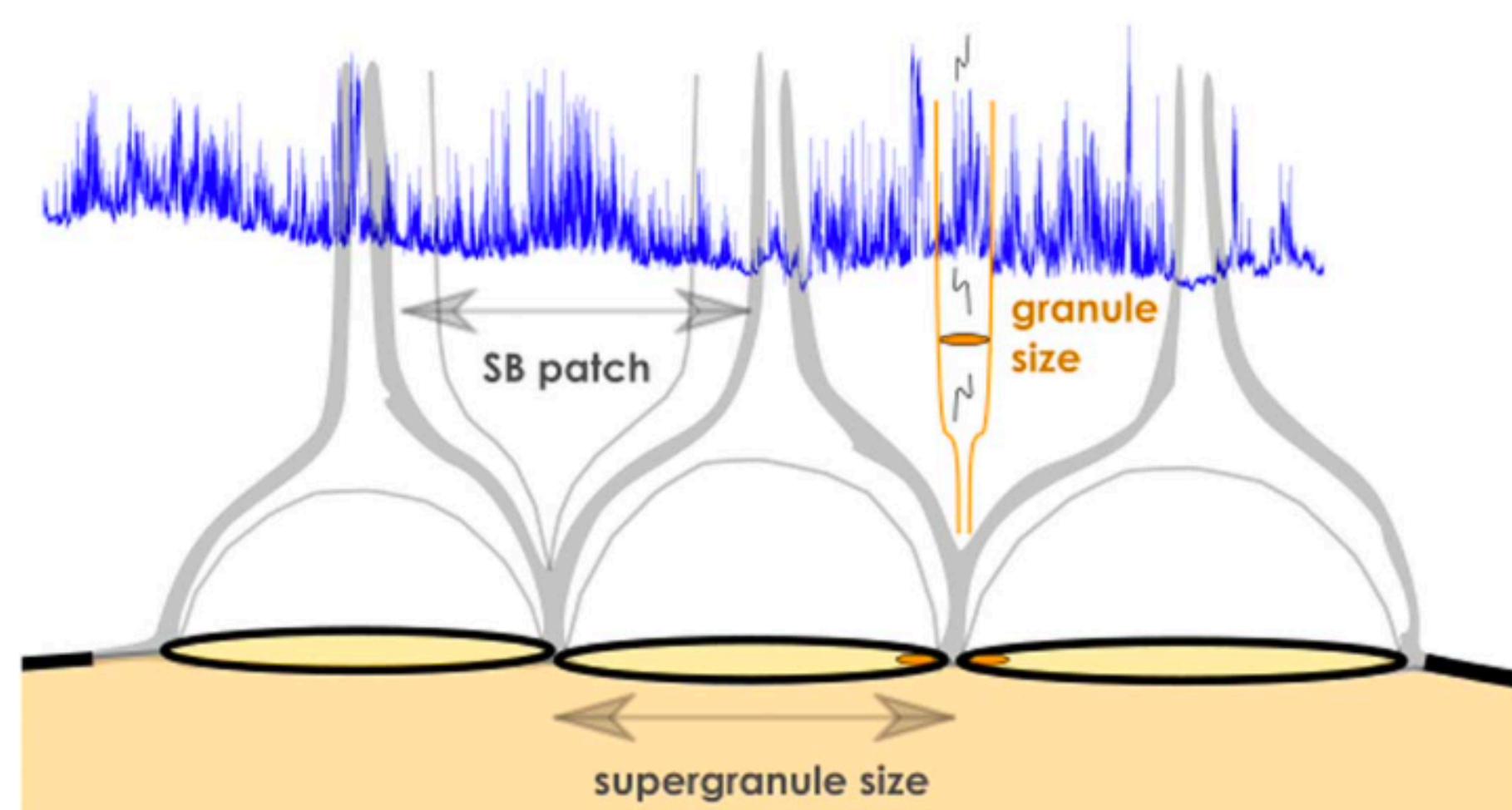


# Switchbacks and surface structures

## Link with the supergranulation

[Farglette et al., 2021]

- Wavelet analysis of SB shows that packets are related to granulation and super granulation scales



# Alfvén wave driven stellar winds

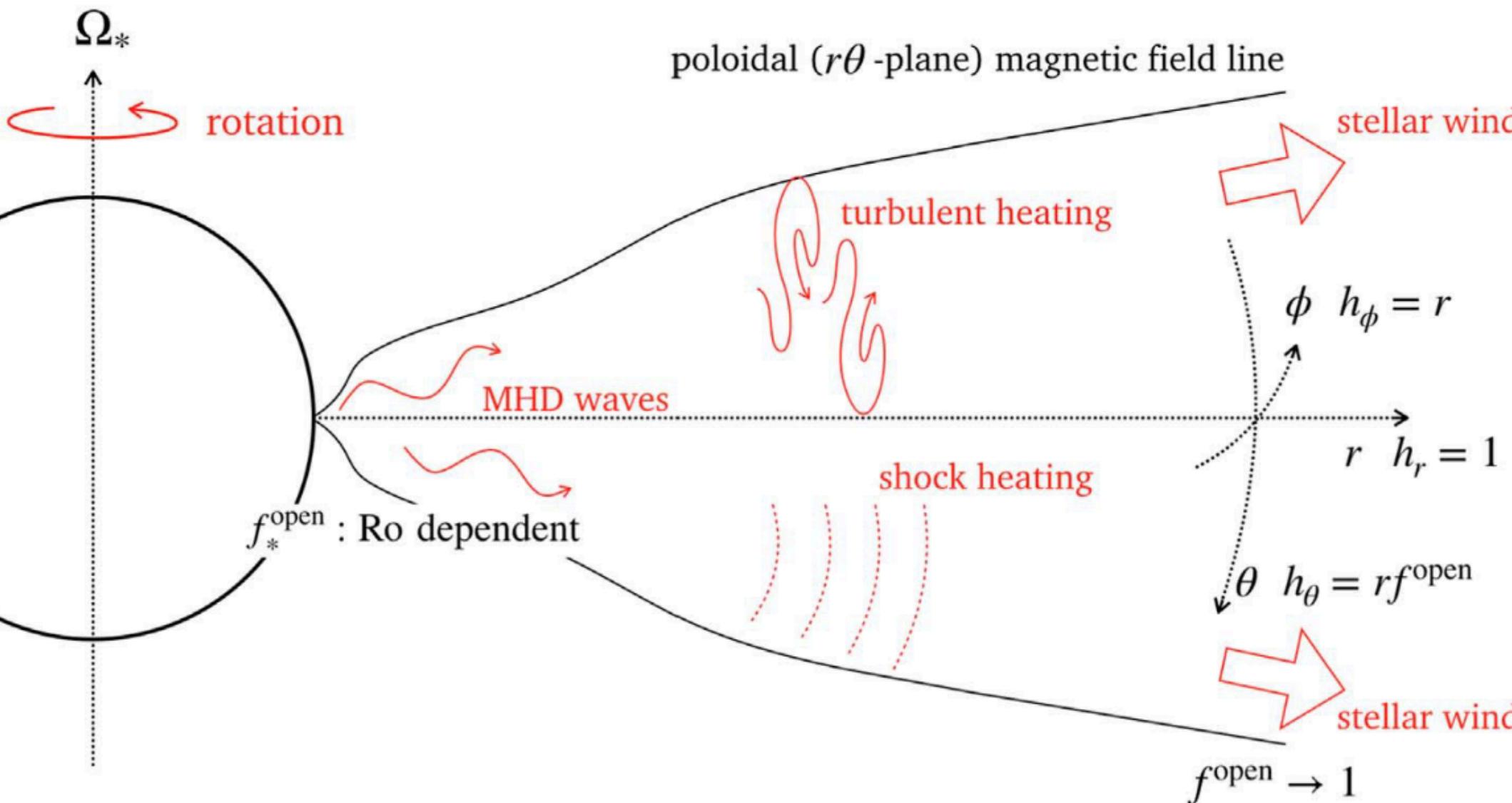
$$\mathcal{P}_{\text{ram}} = \rho v^2$$

## 1D models of fast rotators

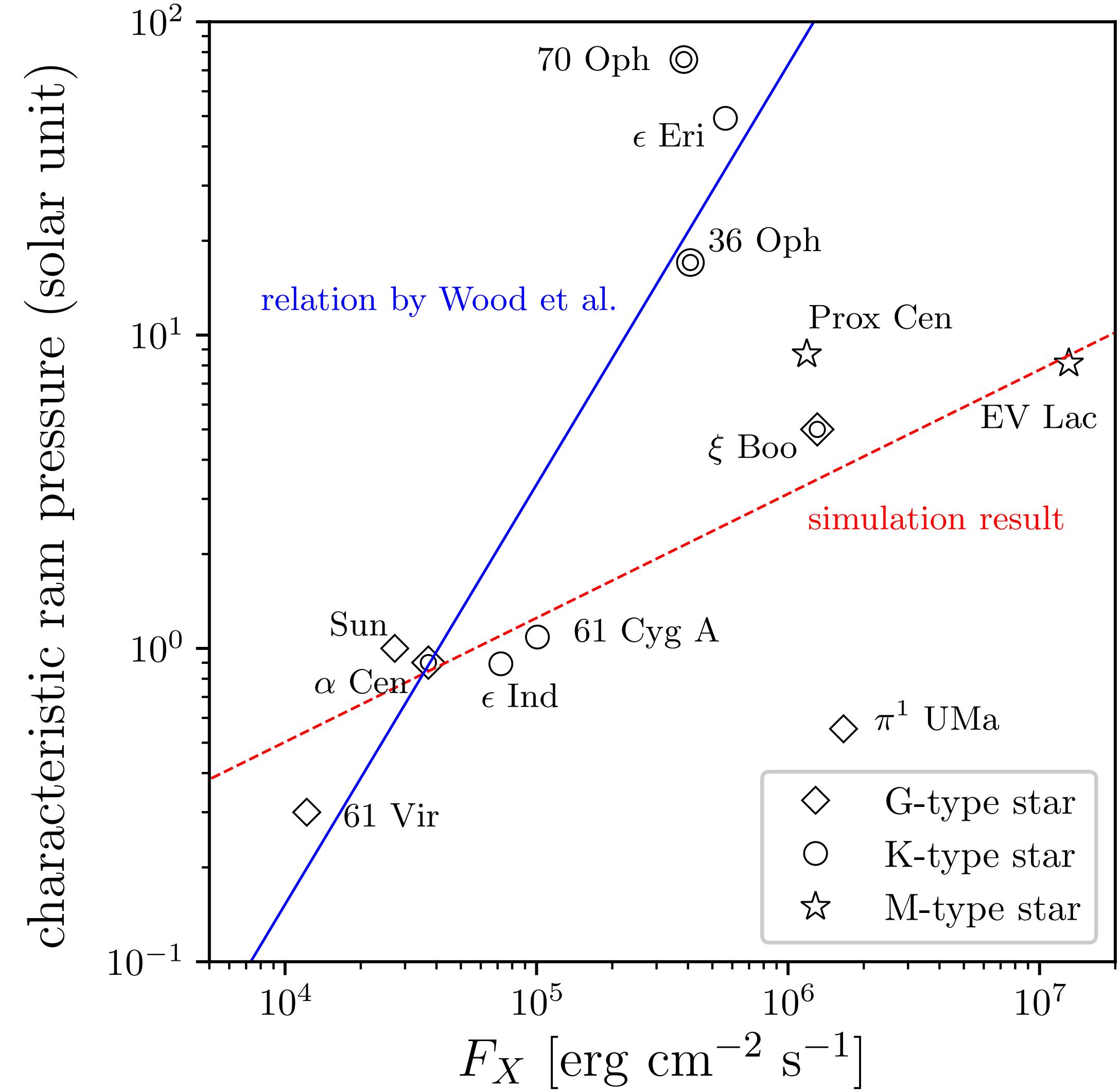
- An equipartition field at the photosphere and the same wave amplitude for all rotation.
- We use the relation :

$$\langle B \rangle \propto P_{\text{rot}}^{-1.2}$$

- We change the ‘expansion’ of the flux tube



[Shoda, Suzuki, ..., Réville et al., 2020]



# Alfvén wave driven stellar winds

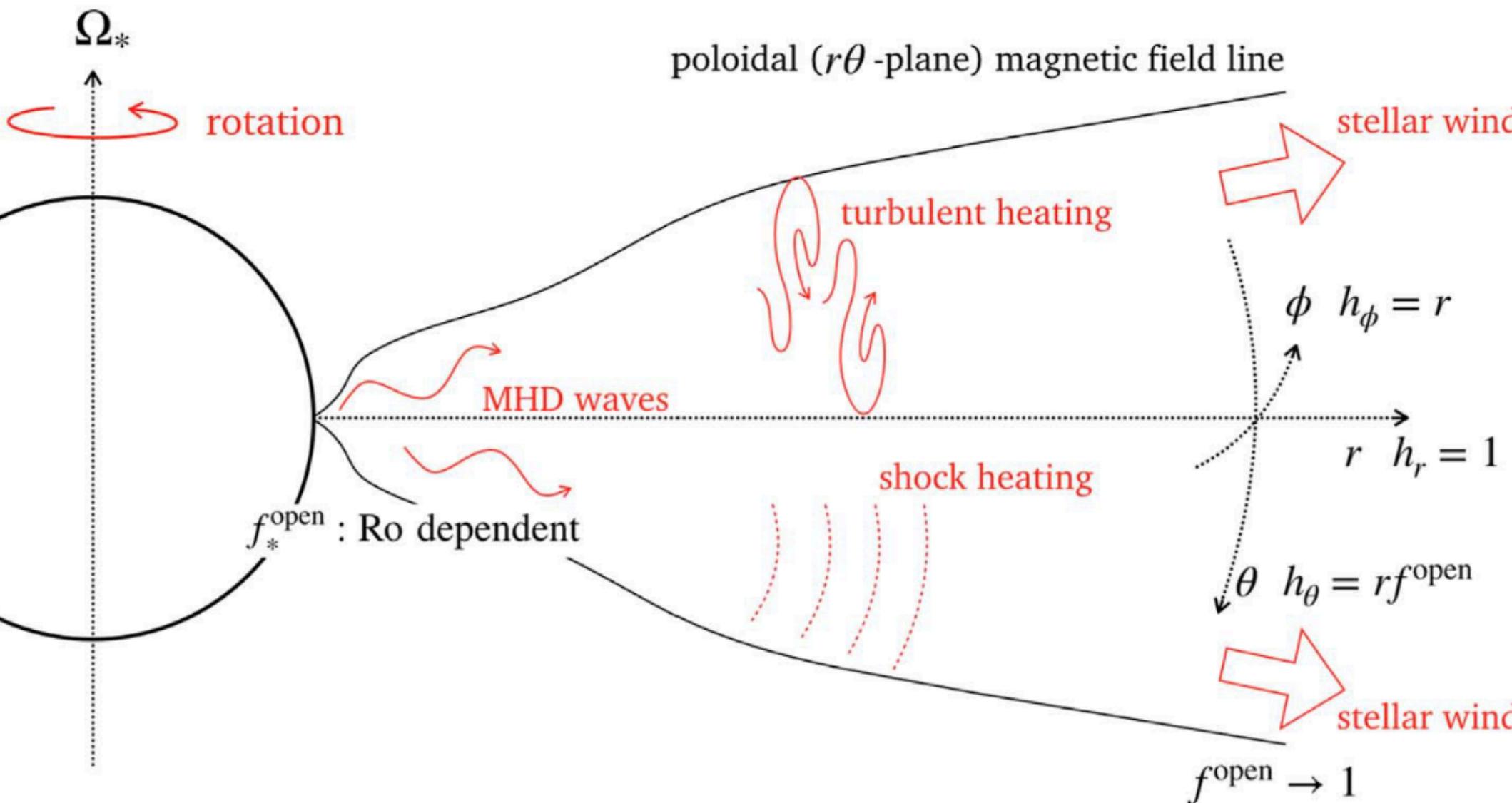
$$\mathcal{P}_{\text{ram}} = \rho v^2$$

## 1D models of fast rotators

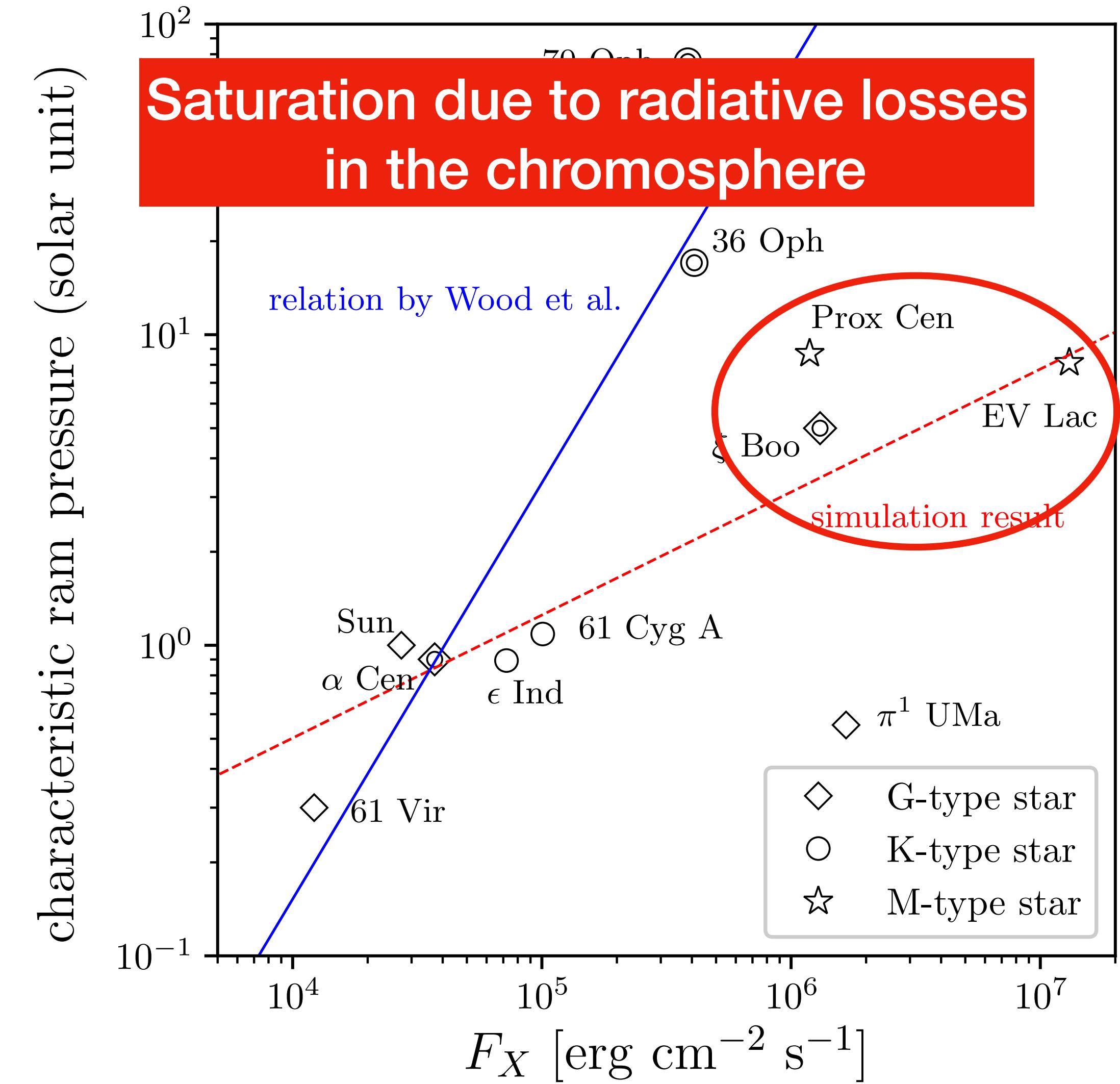
- An equipartition field at the photosphere and the same wave amplitude for all rotation.
- We use the relation :

$$\langle B \rangle \propto P_{\text{rot}}^{-1.2}$$

- We change the ‘expansion’ of the flux tube



[Shoda, Suzuki, ..., Réville et al., 2020]



# Alfvén wave driven stellar winds

## Fully convective stars / M-dwarf

- TRAPPIST-1 System (M-dwarf+ 7 planets)

$$\langle B \rangle = 600 \text{ G} \quad P_{\text{rot}} = 3.3 \text{ d}$$

- Garraffo et al. 2017 (AWSOM)  $1 \dot{M}_{\odot}$

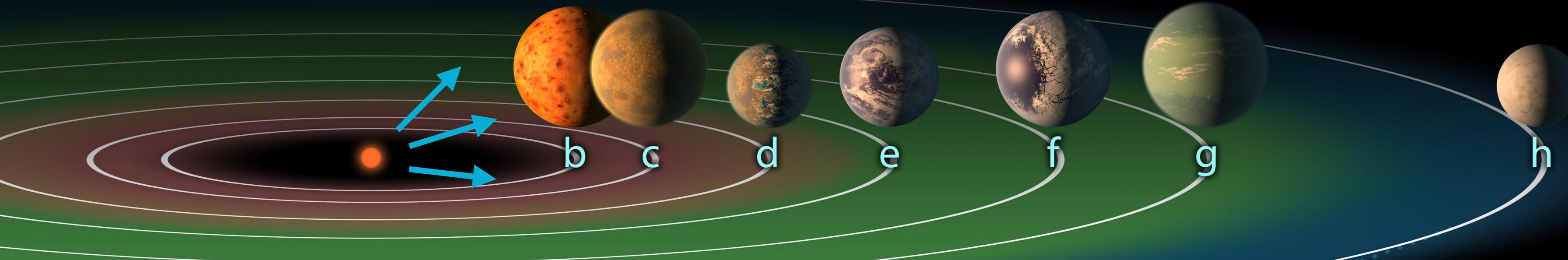
- Dong et al. 2018 (AWSOM)  $0.1 \dot{M}_{\odot}$

*Poynting Flux*

*w/ AW turbulence wind*

$$\dot{M} \propto F_p \propto \rho_{\star} v_{A,\star} \delta v^2$$

TRAPPIST-1 System

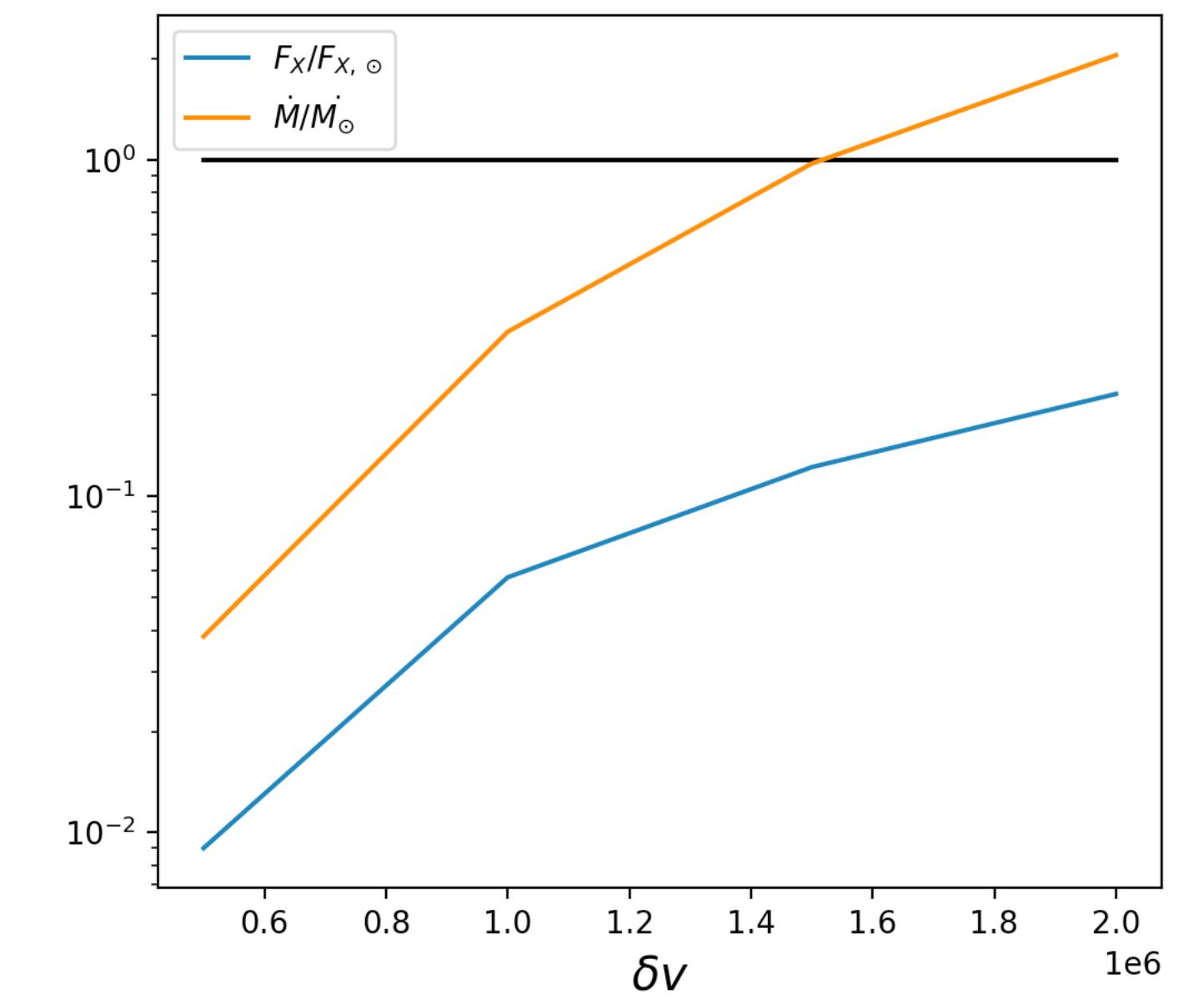
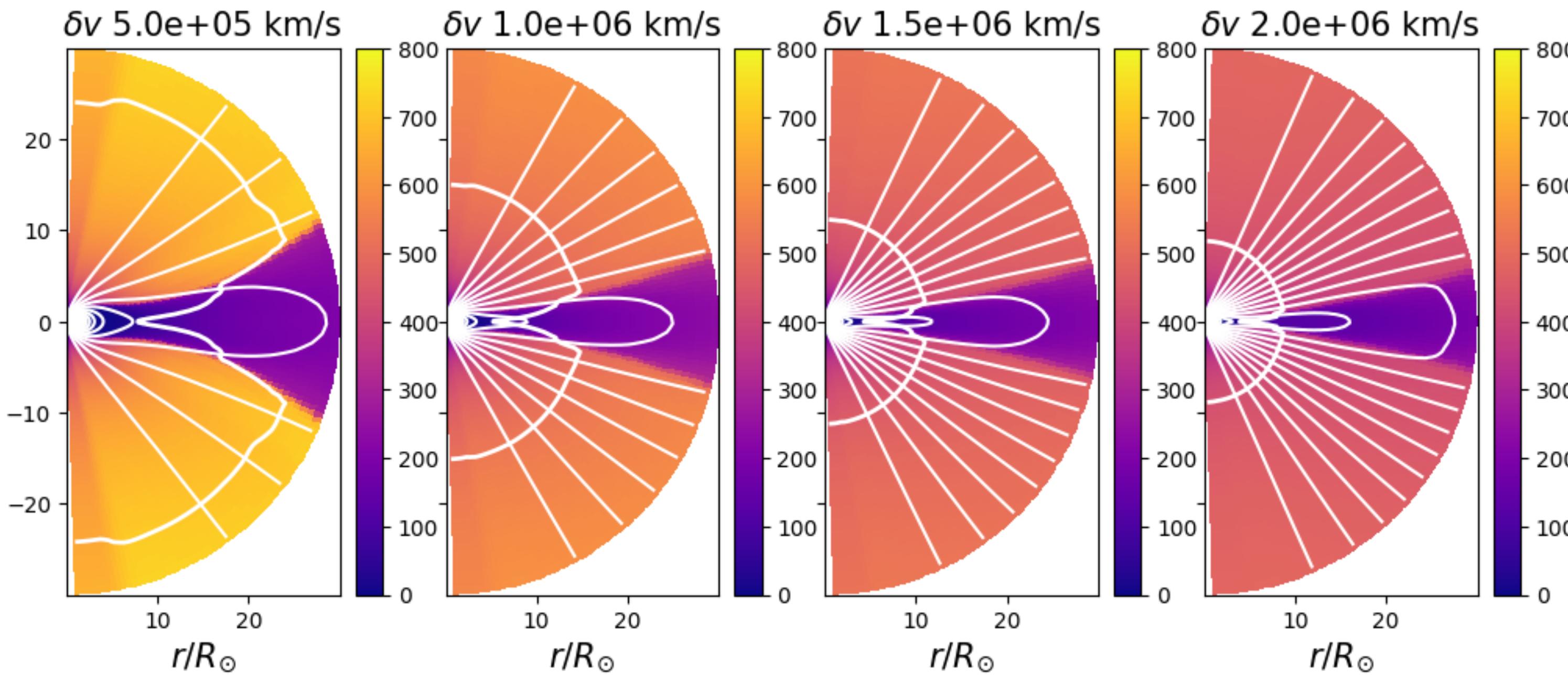


# M-dwarfs stellar winds?

## Estimating the mass loss using X-ray constraints

$$B_\star = 600G$$

[Réville et al., in prep]



- Using Chianti, we integrate the response of the coronal for 2.5D simulation of typical M-dwarf:

*Still below observations -> more small scale structures ?*

# Summary

- *Stellar winds are ubiquitous.*
- *AW turbulence driven models are very efficient at reproducing the solar wind.*
- *Stellar observables, and in particular saturation phases remain mysterious and start to be investigated in the framework of AW turbulence.*
- *Results from the Parker Solar Probe mission show promising leads to bridge the gap! Solar Orbiter will ideally complete the picture.*

**Thanks!**